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DYNAMICS OF PITCHING CONTROL MECHANISM

OF SMALL PROPELLER-TYPE WIND MACHINES

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ABSTRACT

The paper is aiming at exploring the area of pitch control type of regulation used in small propeller-type wind machines. To cover the lack of published information about the real design criteria of such type of regulation, this work is carried out to develop the pertaining theoretical database. The paper investigates the dynamic performance of the centrifugally-activated pitch control mechanisms widely-used in wind machines. A typical pitching mechanism is chosen to be the subject of analysis; namely the crank-slider-crank mechanism which is centrifugally-activated by hub-mounted fly weights. The kinematic and dynamic characteristics of the candidate system are fully analysed. The results are presented in a group of charts in terms of the system dimensionless design-parameters, best suiting the engineers for design purposes.

1. INTRODUCTION

. Pitching is that form of action that causes a change of the setting angle of the propeller blade (or of part of it). As a consequence of that, the angle of attack of each airfoil section along the blade, changes with respect to the relative wind speed. This change has a direct influence on the speed of rotation and/or the power extracted from the wind. The propeller-type small wind energy conversion systems (PT-SWECS) considered in this paper, are normally used for electricity generation. Depending upon the specific application, pitch control in PT-SWECS is used to perform one or more of the following requirements [1-5]:

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- o Self starting of the wind machine.
- o Optimum matching between the wind rotor and the generator used.
- o Limiting the wind machine power-output in the range of operation between the rated wind speed and the cut-out wind speed.
- o Keeping constant rotational speed through the whole operation range of the PT-SWECS while maintaining an optimum output for any particular tip speed ratio.

.Reviewing the commercially-available PT-SWECS, [6], it has been found that blade pitching is a practically proved concept for control. Blade pitching is adopted for control in more than 70 % of the currently produced types of PT-SWECS, [5]&[6]. Furthermore, in all of these machines, the pitchchange mechanism is centrifugally-activated via the centrifugal action of masses rotating with the rotor. The objective of this paper is to study and analyse the so-called hub-mounted fly-weight pitch control mechanism. The paper lays down the governing relations required for designing and/or predicting the performance of the mechanism.

2. KINEMATICS OF THE SYSTEM

The investigated pitch-control mechanism is a typical one which comprises three main components, as shown in Fig.(1). These are: the speed sensor, the output transmitter, and the transmission mechanism. The system contains as much of this mechanism, as the number of the rotor blades.

The pitching mechanism is equivalent to the planar crank-slider-crank mechanism shown in Fig.(2). When the mass is affected by the centrifugal force, it is forced to move in an arc of a circle having the center at point P. Crank 1 will follow the mass movement. Consequently, crank 2 will be forced to move about the hinge P , and point P will move in a circular path of center at P. The slider will slide on crank 1 upward/downward, depending on the direction of motion. The slider translatory motion of Fig.(2), is equivalent to the motion of the slot of the mass lever relative to the pin of the output transmitter of Fig.(1). More details on equivalent mechainisms are provided in [7].

2.1 GEOMETRY OF MOTION

Referring to Fig.(3), and assuming that the mass lever is so designed with 90° angle between its legs, the following relations can be deduced, [6],:

$$\mathbb{X}_{1} = \sqrt{a^{2} + R^{2} + 2 a R \sin \psi}$$

(1)



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Fig.(1): Schematic Diagram of the Investigated Hub-Mounted Fly-Weight Pitching Mechanism



Fig. (2): The Planer Crank Slider Crank

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Fig.(3): Geometrical Notations of The Mechanism

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$I_3 = c \ (\sin \theta \ / \ \sin \psi)$		(2)
c. $sin(\psi + \theta) = R \cdot sin\psi$		(3)
$s = X_3 - R + c$		(4)

The above relations can be expressed, more conveniently, in dimensionless form by introducing the geometrical parameters A & C, where A is the fly-weight arm ratio and C is the blade driving-arm ratio; defined by:

$$A = a/R$$
(5)

$$C = c/R \tag{6}$$

Equations (1) to (4) become:

$$\widetilde{\mathbf{X}}_{1} = \sqrt{1 + A^{2} + 2 A \sin \psi}$$

$$\widetilde{\mathbf{X}}_{3} = C \sin \theta / \sin \psi$$
(7)
(8)

$$\sin \Psi = C \sin(\Psi + \Theta) \tag{9}$$

$$S = X_3 + C - 1$$
 (10)

where:

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$$\bar{X}_1 = X_1/R, \ \bar{X}_3 = X_3/R, \text{ and } S = s/R$$
 (11)

2.2 CRITERION OF MECHANISM INTEGRITY

Fig.(2) shows that crank 1 can drive crank 2 until they are at right angle. This condition occurs on both sides of PP. This is a criterion that puts a limit on the maximum poss² ible movement of the fly-weight mass with the mechanism still having its integrity. At the limiting positions, the following relation is valid:

$$\sin \Psi = + C$$

(12)

(13)

: It follows that the criterion of system integrity is:

 $-C \ge \sin \Psi \ge + C$

The relations defining the geometry of motion of the mechanism, are presented in the form of dimensionless-parameters charts on Figs.(4) to (7). It is evident that for small values of γ , the mechanism movement is approximately linear.









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against Input Angular Position ${\cal V}$

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3. DYNAMICS OF THE SYSTEM

As a control system, the pitching mechanism receives one input which is the rotational speed Ω . The system single output is the setting angle Θ . As the rotor aerodynamic forces consequently change, they are interpretted into a variation of the rotational speed and/or its tendency of variation. This closed loop process continues until a condition of equilibrium, depending upon the operation conditions and the inherent characteristics of the system, is attained. It is worth mentioning that angle θ is not directly the pitch angle. However, it is in direct proportionality to the pitch angle, in a way depends on the transmission mechanism used.

The system under investigation, is highly non-linear due to: its geometry, the centrifugal force being proportional to the square of the rotational speed, the backlash in gears, spring hysterisis, ... etc. Nevertheless, it is generally enough for practical application to limit the degree of non-linearity, [8]. This is achieved by introducing the following assumptions:

o The activated mass is concentrated at its C.G.

- o The spring force-deflection relation is linear.
- o The line of action of the spring force is always radial. o The damping in the system is of the viscous type, i.e., it is proportional to the angular speed θ^* .

o The blade pitching-moment is transferred as a balancing moment M about the axis of rotation where:

$$M = N_{\Omega}^{2} R^{2} (1 + \frac{1}{(TSR)^{2}})$$
(14)

where N is a constant depending mainly on the aerodynamic characteristics of the blade, and TSR is the rotor tip speed ratio. Since TSR is relatively high in the operation range of PT-SWECS, typically between 5 and 15, Eq.(14) can be approximated to:

 $M \cong N n^2 R^2$

(15)

3.1 THE GOVERNING EQUATION

The equation governing the motion of the system can be obtained by balancing the acting forces, moments and inertia terms. The details of this step are given in [6]. The arrived-at equation is introduced hereinafter in its full form. The definition of the variables and constants involved in the following system-dynamics equation, is given in the Nomenclature at the end of the paper.

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 $(C/\overline{X}_3)\cos(\Psi+\Theta)\left[(J/mR^2)(\overline{X}_3^2/C^2)-1\right]\Theta^{**}+\left[(1-C^2)\sin(\Psi/\overline{X}_3^2\right]\Theta^{*2}$ + $\left[(2/ma^2)(\overline{x}_3/C) \cos(\gamma + \theta) \right] \theta^{\circ} + (M/ma^2)(\overline{x}_3/C) \cos(\gamma + \theta)$ + $(K/mA)(1 - \overline{L/X_1}) \cos \psi = \left[(A + \sin \psi)/A \right] \Omega^{\circ} + (\cos \psi/A) \Omega^2$ (16)

The above equation is non-linear, second-order, second-degree ordinary differential equation with time-dependent coefficie-.nts. The input function Ω appears in the governing equation in two forms; its square(Ω) and its rate of change (Ω). Thus the system is expected to respond either to the rotational speed change or to its tendancy to change, or to both of them. This is a favourable characteristics, [8] to [10]. In what follows, the dynamic analysis of the system is performed via an exact steady state analysis only. A transient-response analysis, not included in the present paper because of the limited space, is provided in [6].

3.2 EXACT STEADY STATE ANALYSIS

The steady state, which is alternatively called the equilibrium running state, is obtained, as described in [11], by setting θ° , θ° and Ω° equal to zero in Eq.(16). Doing so, the following expression, defining the state of steady running of the system, is obtained:

$$\overline{D} = \sqrt{\frac{(1 - \overline{L}/\overline{X}_{1})}{\left[1 - H_{1}(\overline{X}_{3}/AC) \left[\cos(\psi + \theta)/\cos\psi\right]\right]}}$$
(17)

By examining the above equation, and in the light of that for practical application, L is generally less than unity and X is more than unity, it can be seen that the system cannot physically exist unless the following expression is satisfied:

$$H_{1} \langle (AC/\bar{X}_{3}) \left[\cos \psi / \cos(\psi + \theta) \right]$$
(18)

This criterion puts an upper limit to the design parameter H₁ (i.e., the ratio N/m) which must not be exceeded at any operating point of the mechanism defined by angle Ψ and its geometry. The system existence criterion defined by the inequality in Eq.(18), is represented in parametric non-dimensional chart in Fig.(8). It is readily apparent that the above criterion is a direct correlation between the blade characteristics defined by N and the system activation element defined by the mass m. The system can be reduced to a simple centrifugal speed sensing device by setting H₁=0. As it will be shown later, the ratio H₁ is almost the major factor affecting the system

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Going further, the system steady state running characteristics are fully defined, [11] & [12], by defining: the input/ output relation $\Omega - \Theta$ and the input/gain relation $\Omega - \overline{G}$. The first one is defined by Eq.(17), while the second is simply obtained by differentiating Eq.(17), giving:

$$f = d\theta/d \Lambda$$
 (19)

$$\frac{1}{\kappa_{1}\bar{\kappa}[\kappa_{7}(1-\bar{\kappa}^{2})-(\bar{L}/\bar{x}_{1})(\kappa_{7}-2\Lambda\cos^{2}\psi/\bar{x}_{1}^{2})]}$$
(20)

where:

$$K_{1} = (C/\overline{X}_{3}) \cos(\psi + \theta)$$
(21)

$$K_7 = (\sin 2\Psi / \bar{X}_3 K_1^2) - 2 \sin \Psi$$
 (22)

Figs.(9) and (10) are examples for the steady state characteristics defined by Eqs.(17) and (20), respectively. The tremendous effect of the design parameter H_1 on the system behaviour, is quite clear in the two figures.

Depending upon the values of A, C, and L, the characteristic line for a certain value of H_1 takes one of the shapes shown on Fig.(11) or Fig.(12). The operating line shown in Fig.(11) is a continuous function starts at P_1 and ends at P_2 after reaching a peak where $\bar{A} = \bar{A}_1$ and $\theta^1 = \theta(\bar{A}_1)$. Fig.(12) shows the case when the operating line is not a continuous function, and where a phenomenon of branching occurs at the point P_3 where $\bar{A} = 0$ and $\theta = \theta_B$.

Since the negative gain regimes imply some kind of instability, [9] & [10], it is recommended to limit the operating regime of the system to be within the positive gain region. In the case of Fig.(11), the operating regime is defined by: $\theta \leq \theta(\bar{n}_{max})$, whilst in Fig.(12) it is defined by $\theta_{\rm B} \leq \theta \leq \theta(\bar{n}_{max})$.

Since the branching phenomenon is not a desirable one, it is also recommended to keep a certain correlation between the parameters A, C, and L, to avoid the occurrence of this phen-: omenon. To define this correlation, we proceed as follows:

The branching point is always at $\bar{\Lambda} = 0$. At this point, $\theta = \theta_{\rm B}$ where $\theta_{\rm B} \ge$ the minimum physical θ . From Eq.(17), $\bar{\Lambda} = 0$ when $1 - \bar{L}/\bar{X}_1 = 0$, which means (using Eq.(7)) that $\sin \psi(\bar{\Lambda} = 0) =$ $(\bar{L}^2 - A^2 - 1)/2A$. Provided that $\psi(\bar{\Lambda} = 0) \ge \psi_{\rm min}$ and since $\sin \psi_{\rm min} = -C$, then the branching phenomenon will occur if $\bar{L}^2 \ge 1 + A^2 - 2AC$. This means that the necessary and suffic-

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Fig.(10): Typical Example of the System Equilibrium-Running Characteristics (Input/ Static-Gain Variation)



Fig.(11): Typical Shape of the System Input/Output Characteristics at Equilibrium



Fig.(12): The Branching Phenomenon

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ient condition to avoid the branching phenomenon is given by: $\overline{L} < 1 + A^2 - 2AC$ (23)

The chart shown in Fig.(13) is prepared on the basis of the above correlation to guide the proper choice of A, C, and L. The other border of the positive gain regime is defined by $\overline{\Omega} = \overline{\Omega}$. The locus of $\overline{\Omega}$ can be found from Eq.(20) by setting $\overline{G} = \infty$, which yields:

$$\bar{n}_{\max} = \left[1 - (\bar{L}/\bar{X}_{1}) (1 - \frac{2 \Lambda \cos^{2} \gamma}{K_{7} \bar{X}_{1}^{2}}) \right]$$
(24)

Fig.(14) presents a typical example of such locus.

4. CONCLUSIONS

The design and performance criteria of a typical pitch-control mechanism, have been studied. The kinematic and dynamic characteristics of the hub-mounted fly-weights, centrifugally activated, crank-slider-crank mechanism, are fully analysed. The analysis provided in this paper constitutes an adequate base for the designer to choose the system parameters. Depending upon the aimed-at performance of the system, and using the charts and information provided by this paper, it is possible to design and synthesize a system that will perform : as required . Moreover, the performance of an existing system can be improved statically and dynamically.

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6. NOMENCLATURE

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fly-weight arm ratio; Eq.(5)
 A
      length of fly-weight arm; Fig.(3)
 a
 C
      moment-arm ratio; Eq.(6)
      arm of driving output moment; Fig.(3)
 С
 G
      dimensionless static gain of system; G=G/\sqrt{K/m}, where G
      is the static gain
 H<sub>J</sub>1
      pitching-moment ratio; H<sub>1</sub>=N/m
      equivalent polar moment of inertia of driven masses
 Κ
      spring force/deflection constant of proportionality
 KI
KI
T7
      system parameter; Eq.(21)
      system parameter; Eq.(22)
      spring-setting ratio; L=L/R, where L is component of
      spring-setting in the radial direction
- m
      mass of the whole centrifugally-activated fly-weight
 M
      equivalent blade pitching-moment transferred to the
      rotor axis; Eqs.(14) & (15)
 N
      pitching-moment constant of proportionality; Eq.(14)
 PT-SWECS
              propeller-type small wind energy conversion system
 R
      characteristic hub-radius of mechanism; Fig.(3)
 S
      slot-length ratio; Eq.(11)
      slot length; Figs.(2) & (3)
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 TSR tip speed ratio of rotor
 radial distance of fly-weight; Fig.(3)
      dimensionless X<sub>1</sub>; Eq.(11)
      distance between pivot P and contact point P, on slot;
      Figs.(2) & (3)
 I
X
X
Y
     dimensionless X<sub>3</sub>; Eq.(11)
system viscous-damping coefficient
      angular position of fly-weight arm; Fig.(3)
 θ
     angular-position output of mechanism; Fig.(3)
 {\boldsymbol{\theta}}_{\boldsymbol{B}}^{i}
     value of \theta at branching point; Fig.(12)
     rate of change of \Theta with time
 0°°
     rate of change of \theta^* with time
S
     rotor speed of rotation; Fig.(1)
ഹ്
     rate of change of \Omega with time
 Ā
      dimensionless \Omega; \overline{\Lambda} = \Lambda / \sqrt{K/m}
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