

EXACT SOLUTION OF BISWAS-MILOVIC EQUATION USING NEW EFFICIENT METHOD

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ABSTRACT. Biswas and Milovic proposed a generalized model for the NLSE that accounts for several imperfections in the fiber during long distance transmission of these pulses. This paper studies the exact solution of the Biswas-Milovic equation with Kerr law nonlinearity by a new efficient method. The results reveal that the method is explicit, effective, and easy to use.

1. INTRODUCTION

Biswas-Milovic equation is a generalized version of the familiar nonlinear Schrödinger's equation describing the propagation of solitons through optical fibers for trans-continental and trans-oceanic distances. A broad class of analytical solution methods and numerical solution methods were used to handle these problems [1-23]. The nonlinear Schrödinger's equation is studied in various areas of Applied Mathematics, Theoretical Physics and Engineering. In particular, it appears in the study of Nonlinear Optics, Plasma Physics, Fluid Dynamics, Biochemistry and many other areas [1-43].

The Biswas-Milovic equation is given by [22, 23]:

$$i(q^m)_t + a(q^m)_{xx} + bF(|q|^2)q^m = 0,$$

where q is a complex valued function, while x and t are the two independent variables. The coefficient a and b are constants where $ab > 0$, and parameter $1 \leq m < 2$.

In this paper we outline a reliable strategy of the new homotopy perturbation method for solving the case $m = 1$ of Biswas Milovic equation

$$\begin{aligned} iq_t + aq_{xx} + bF(|q|^2)q &= 0, \\ q(x, 0) &= g(x). \end{aligned} \tag{1}$$

The Biswas-Milovic equation plays an important role in mathematical physics. It is very well known that the dynamics of these solitons propagate through these fibers for long distances in a matter of a few femto-seconds. This dynamics is governed by the nonlinear Schrödinger's equation (NLSE). The optical fiber industry

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has grown profoundly based on this technology. Recently, the homotopy perturbation method (HPM) introduced by He [24, 25] has been widely used for solving various integral equations arising from real world modelling, for example thin film flow, heat transfer, and many others [26-37]. The idea behind this method is that the solution is considered as the sum of an infinite series, which converges rapidly to the exact solution. In this paper we present a modified version of HPM and call it NHPM, which performs much better than the HPM. The new homotopy perturbation method (NHPM) was applied to linear and nonlinear PDEs and hyperbolic Telegraph equation. The rest of the paper is organized as follows. In section 2 we apply the method on the non-linear Biswas-Milovic equation. The results of numerical experiments are presented in section 3. Section 4 is dedicated to a brief conclusion. Finally some references are introduced at the end. The Biswas-Milovic equation plays an important role in mathematical physics. It is very well known that the dynamics of these solitons propagate through these fibers for long distances in a matter of a few femto-seconds. This dynamics is governed by the nonlinear Schrödinger's equation (NLSE). The optical fiber industry has grown profoundly based on this technology.

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2. APPLICATION OF NHPM TO BISWAS-MILOVIC EQUATION

For the purpose of applications illustration of the methodology of the proposed method, using new homotopy perturbation method, we consider the following differential equation

$$A(u(X, t) - f(r)) = 0, \quad r \in \Omega, \quad (2)$$

$$B(u(X, t), \partial u / \partial n) = 0, \quad r \in \Gamma, \quad (3)$$

where A is a general differential operator, $f(r)$ is a known analytic function, B is a boundary condition, Γ is the boundary of the domain Ω , and $X = (x_1, x_2, \dots, x_n)$. The operator A can be generally divided into two operators, L and N , where L is a linear, while N is a non-linear operator. Equation (2) can be written as follows:

$$L(u) + N(u) - f(r) = 0. \quad (4)$$

Using the homotopy technique, we construct a homotopy $U(r, p) : \Omega \times [0, 1] \rightarrow \mathbb{R}$ which satisfies:

$$H(U, p) = (1 - p)[L(U) - L(u_0)] + p[A(U) - f(r)] = 0, \quad p \in [0, 1], \quad r \in \Omega, \quad (5)$$

or

$$H(U, p) = L(U) - L(u_0) + pL(u_0) + p[N(U) - f(r)] = 0, \quad (6)$$

where $p \in [0, 1]$, is called homotopy parameter, and u_0 is an initial approximation for the solution of Eq. (1), which satisfies the boundary conditions. Obviously from Eqs. (5) and (6) we will have

$$H(U, 0) = L(U) - L(u_0) = 0, \quad (7)$$

$$H(U, 1) = A(U) - f(r) = 0. \quad (8)$$

We can assume that the solution of (5) or (6) can be expressed as a series in p , as follows:

$$U = U_0 + pU_1 + p^2U_2 + \dots \quad (9)$$

Setting $p = 1$, produces the approximate solution of Eq. (2), which could be written in the following form:

$$u = \lim_{p \rightarrow 1} U = U_0 + U_1 + U_2 + \dots$$

Now let us write the Eq. (6) in the following form

$$L(U(X, t)) = u_0(X, t) + p[f(r(X, t)) - u_0(X, t) - N(U(X, t))]. \quad (10)$$

By applying the inverse operator, L^{-1} to both sides of Eq. (10), we have

$$U(X, t) = L^{-1}(u_0(X, t)) + p \left(L^{-1}(f(r)) - L^{-1}(u_0(X, t)) - L^{-1}(N(U(X, t))) \right). \quad (11)$$

Suppose that the initial approximation of Eq. (2) has the form

$$u_0(X, t) = \sum_{n=0}^{\infty} a_n(X)P_n(t), \quad (12)$$

where $a_1(X), a_2(X), a_3(X), \dots$ are unknown coefficients and $P_0(t), P_1(t), P_2(t), \dots$ are specific functions depending on the problem. Now by substituting (9) and (12) into the Eq. (11), we get

$$\begin{aligned} \sum_{n=0}^{\infty} p^n U_n(X, t) = U(X, t) = L^{-1} \left(\sum_{n=0}^{\infty} a_n(X)P_n(t) \right) \\ + p \left(L^{-1}(f(r)) - L^{-1} \left(\sum_{n=0}^{\infty} a_n(X)P_n(t) \right) - L^{-1} \left(N \left(\sum_{n=0}^{\infty} p^n U_n(X, t) \right) \right) \right). \end{aligned} \quad (13)$$

Comparing coefficients of terms with the identical powers of p , leads to

$$\begin{aligned}
 p^0 : U_0(X, t) &= L^{-1}(\sum_{n=0}^{\infty} a_n(X)P_n(t)), \\
 p^1 : U_1(X, t) &= L^{-1}(f(r)) - L^{-1}(\sum_{n=0}^{\infty} a_n(X)P_n(t)) - L^{-1}N(U_0(X, t)), \\
 p^2 : U_2(X, t) &= -L^{-1}N(U_0(X, t), U_1(X, t)), \\
 p^3 : U_3(X, t) &= -L^{-1}N(U_0(X, t), U_1(X, t), U_2(X, t)), \\
 &\vdots \\
 p^j : U_j(X, t) &= -L^{-1}N(U_0(X, t), U_1(X, t), U_2(X, t), \dots, U_{j-1}(X, t)), \\
 &\vdots
 \end{aligned} \tag{14}$$

Now if we solve these equations in such a way that $U_1(X, t) = 0$, then Eqs. (14) results in $U_1(X, t) = U_2(X, t) = \dots = 0$.

Therefore the exact solution may be obtained as the following.

$$u(X, t) = U_0(X, t) = \sum_{n=0}^{\infty} a_n(X)P_n(t). \tag{15}$$

It is worthwhile to mention that if $f(r)$, and $u_0(X, t)$, are analytic at $t = t_0$, then their Taylor series defined as the following

$$\begin{aligned}
 u_0(X, t) &= \sum_{n=0}^{\infty} a_n(X)(t - t_0)^n, \\
 f(r(X)) &= \sum_{n=0}^{\infty} a_n^*(X)(t - t_0)^n,
 \end{aligned}$$

can be used in Eq. (13), where $a_1(X), a_2(X), a_3(X), \dots$ are unknown coefficients and $a_1^*(X), a_2^*(X), a_3^*(X), \dots$ are known ones, which should be computed.

To show the capability of the method, we apply the NHPM to some examples in the next section.

3. ILLUSTRATIVE EXAMPLES

In this section we present some numerical results of our scheme for the nonlinear Biswas-Milovic [22, 23].

Example 1. Consider the non linear B-M equation

$$iq_t + q_{xx} + 2|q|^2 q = 0, \tag{16}$$

The initial conditions are given by

$$q(x, 0) = e^{ix}.$$

To solve Eq. (16) by the NHPM, we construct the following homotopy:

$$\frac{\partial Q}{\partial t} = q_0 - p \left(q_0 - iQ_{xx} - 2i|Q|^2 Q \right). \tag{17}$$

Applying the inverse operator, $L^{-1} = \int_0^t (\cdot) dt$ to the both sides of Eq. (17), we obtain

$$Q(x, t) = q(x, 0) + \int_0^t q_0 dt - p \int_0^t \left(q_0 - iQ_{xx} - 2i|Q|^2 Q \right) dt, \tag{18}$$

Suppose the solution of Eq. (18) to have the following form

$$Q = Q_0 + pQ_1 + p^2Q_2 + \dots, \tag{19}$$

where Q_i are unknown functions which should be determined. Substituting Eq. (19) into Eq. (18), collecting the same powers of p , and equating each coefficient of p to zero, results in

$$\begin{aligned} p^0 : Q_0(x, t) &= Q(x, 0) + \int_0^t q_0(x, t) dt, \\ p^1 : Q_1(x, t) &= \int_0^t \left(q_0 + iQ_{0xx} + 2i|Q_0|^2 Q_0 \right) dt, \\ p^2 : Q_2(x, t) &= \int_0^t \left(iQ_{1xx} + 2i|Q_0|^2 Q_1 + 4i|Q_0||Q_1|Q_0 \right) dt, \\ &\vdots \end{aligned}$$

Setting $p = 1$, results in the approximate solution of Eq. (19)

$$q = \lim_{p \rightarrow 1} Q_0 + pQ_1 + p^2Q_2 + \dots,$$

Assuming $q_0(x, t) = \sum_{n=0}^{\infty} a_n(x)t^n$, $Q(x, 0) = q(x, 0)$, and solving the above equation for $Q_1(x, t)$ leads to the result

$$Q_1(x, t) = \left(a_0(x) - \frac{1}{2}ie^{ix} \right) t + \left(-\frac{1}{2}a_1(x) - \frac{1}{12}e^{ix} \right) t^2 + \dots$$

Vanishing $Q_1(x, t)$ lets the coefficients $a_n(x)$ ($n = 1, 2, 3, \dots$) to take the following values

$$a_0(x) = \frac{1}{2}ie^{ix}, a_1(x) = -\frac{1}{6}e^{ix}, a_2(x) = -\frac{1}{24}ie^{ix}, \dots$$

Therefore, the exact solutions of the Eq. (16) can be expressed as

$$\begin{aligned} q(x, t) = Q_0(x, t) &= e^{ix} + a_0(x)t + \frac{1}{2}a_1(x)t^2 + \frac{1}{6}a_2(x)t^3 \\ &+ \frac{1}{12}a_3(x)t^4 + \frac{1}{20}a_4(x)t^5 + \dots = e^{i(x+t)}. \end{aligned}$$

Example 2. Consider the non-linear B-M equation

$$iq_t + q_{xx} + |q|^4 q = 0, \quad (20)$$

with the initial condition

$$q(x, 0) = e^{ix}.$$

To solve Eq. (20) by the NHPM, we construct the following homotopy

$$\frac{\partial Q}{\partial t} = q_0 - p \left(q_0 - iQ_{xx} - i|Q|^4 Q \right). \quad (21)$$

Integrating of Eq. (21) leads to

$$Q(x, t) = q(x, 0) + \int_0^t q_0 dt - p \int_0^t \left(q_0 - iQ_{xx} - i|Q|^4 Q \right) dt, \quad (22)$$

Suppose the solutions of equation (19) have the form (19), substituting Eq. (19) into Eq. (22), collecting the terms with the same powers of p , and equating each

coefficient of p to zero, results in

$$\begin{aligned} p^0 : Q_0(x, t) &= Q(x, 0) + \int_0^t q_0(x, t) dt, \\ p^1 : Q_1(x, t) &= \int_0^t \left(q_0 + iQ_{0xx} + i|Q_0|^4 Q_0 \right) dt, \\ p^2 : Q_2(x, t) &= \int_0^t \left(iQ_{1xx} + i|Q_0|^4 Q_1 + 4i|Q_0|^3 |Q_1| Q_0 \right) dt, \\ &\vdots \end{aligned}$$

Assuming $q_0(x, t) = \sum_{n=0}^{\infty} a_n(x)t^n$, $Q(x, 0) = q(x, 0)$, and we set the Taylor series of $Q_1(x, t)$ at $t = 0$, equal to zero, we have

$$Q_1(x, t) = (a_0(x))t + \left(-\frac{1}{2}a_1(x) - \frac{1}{12}a_0(x)e^{ix} \right) t^2 + \dots$$

It can be easily shown that

$$a_0(x) = 0, a_1(x) = 0, a_2(x) = \frac{1}{2}x, a_3(x) = 0, \dots$$

Therefore we gain the solution of Eq. (20) as

$$\begin{aligned} q(x, t) = Q_0(x, t) &= e^{ix} + a_0(x)t + \frac{1}{2}a_1(x)t^2 + \frac{1}{6}a_2(x)t^3 \\ &+ \frac{1}{12}a_3(x)t^4 + \frac{1}{20}a_4(x)t^5 + \dots = e^{ix}. \end{aligned}$$

Example 3. In this example we will apply the new method to handle the BME with parabolic law nonlinearity

$$iq_t + 4q_{xx} + 2(|q|^2 + |q|^4)q = 0, \quad (23)$$

with the initial condition

$$q(x, 0) = e^{-\frac{1}{2}ix}.$$

To solve Eq. (23) by the NHPM, we construct the following homotopy

$$\frac{\partial Q}{\partial t} = q_0 - p \left(q_0 - 4iQ_{xx} - 2i(|Q|^4 + |Q|^2)Q \right). \quad (24)$$

Integrating of Eq. (24) leads to

$$Q(x, t) = q(x, 0) + \int_0^t q_0 dt + p \int_0^t \left(q_0 + 4iQ_{xx} + 2i(|Q|^4 + |Q|^2)Q \right) dt, \quad (25)$$

Suppose the solutions of equation (23) have the form (19), substituting Eq. (19) into Eq. (25), collecting the terms with the same powers of p , and equating each coefficient of p to zero, results in

$$\begin{aligned} p^0 : Q_0(x, t) &= Q(x, 0) + \int_0^t q_0(x, t) dt, \\ p^1 : Q_1(x, t) &= \int_0^t \left(q_0 + 4iQ_{0xx} + 2i(|Q_0|^4 + |Q_0|^2)Q_0 \right) dt, \\ p^2 : Q_2(x, t) &= \int_0^t \left(4iQ_{0xx} + 2i(|Q_0|^4 + |Q_0|^2)Q_1 + 2i(4|Q_0|^3 |Q_1| + 2|Q_0| |Q_1|)Q_0 \right) dt, \\ &\vdots \end{aligned}$$

Assuming $q_0(x, t) = \sum_{n=0}^{\infty} a_n(x)t^n$, $Q(x, 0) = q(x, 0)$, and we set the Taylor series of $Q_1(x, t)$ at $t = 0$, equal to zero, we have

$$Q_1(x, t) = \left(a_0(x) - \frac{1}{2}ie^{-\frac{1}{2}ix} \right) t + \left(-\frac{1}{2}a_1(x) - \frac{1}{12}e^{-\frac{1}{2}ix} \right) t^2 + \dots$$

It can be easily shown that

$$a_0(x) = \frac{1}{2}ie^{-\frac{1}{2}ix}, a_1(x) = -\frac{1}{6}e^{-\frac{1}{2}ix}, a_2(x) = -\frac{1}{24}ie^{-\frac{1}{2}ix}, \dots$$

Therefore we gain the solution of Eq. (23) as

$$q(x, t) = Q_0(x, t) = e^{-\frac{1}{2}ix} + a_0(x)t + \frac{1}{2}a_1(x)t^2 + \frac{1}{6}a_2(x)t^3 + \frac{1}{12}a_3(x)t^4 + \frac{1}{20}a_4(x)t^5 + \dots = e^{-\frac{1}{2}ix+it}.$$

4. CONCLUSION

In this work, we have used a new homotopy perturbation method for finding solutions of the Biswas-Milovic equation. All the examples show that the NHPM is a powerful mathematical tool for solving Biswas-Milovic equation. It is also a promising method for solving other nonlinear equations. It should be mentioned that we have successfully applied the NHPM to a wide variety of initial-type differential equations of hyperbolic and system of differential equations and we obtained the exact solutions of these systems. Therefore, we believe that new method is a promising technique in finding the exact solutions for a wide variety of mathematical problems. The computations associated with the examples in this article were performed using Maple 13.

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